

DT INVARIANTS AND A
NON-PERTURBATIVE TOPOLOGICAL
STRING PARTITION FUNCTION

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1. Overview

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 - ▶ Flexible enough to include the known dualities.
 - ▶ Sufficiently precise to include the concrete predictions.

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- They arise naturally in string theory as categories of branes in topological twists (A- and B-model).
- Equivalences of such categories provide a language for describing physical dualities (e.g. HMS).
- Want to build a mathematical understanding of string theory based on such categories.

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 - ▶ a group homomorphism $Z: K_0(\mathcal{D}) \rightarrow \mathbb{C}$,together satisfying some axioms.
- Heuristic example: suppose $X = X(I, \omega)$ is a compact Calabi-Yau threefold, and $\mathcal{D} = \mathcal{D}^b \text{Fuk}(X, \omega)$. Then there should be a stability condition $\sigma(I) = (Z, \mathcal{P})$ on \mathcal{D} with
 - ▶ $\mathcal{P}(\phi) = \{\text{special Lagrangians of phase } \phi\} \subset \mathcal{D}$,
 - ▶ $Z(L) = \int_L \Omega^{3,0} \in \mathbb{C}$.

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- In recent work with Ian Strachan we explained that this structure should be a complex hyperkähler structure on the total space of the tangent bundle.
- We can try to construct it by solving a class of Riemann-Hilbert problems in the complex plane with discontinuities prescribed by the DT invariants.

2. Riemann-Hilbert problems from DT theory.

See: "Riemann-Hilbert problems from Donaldson-Thomas theory", arxiv.

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Introduce the algebraic torus

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Assume now that \mathcal{D} has the CY_3 property. Then the Euler form

$$\langle [E], [F] \rangle = \sum_{i \in \mathbb{Z}} (-1)^i \dim_{\mathbb{C}} \text{Hom}_{\mathcal{D}}^i(E, F[i]): \Gamma \times \Gamma \rightarrow \mathbb{Z}$$

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is skew-symmetric. This gives an invariant Poisson structure on \mathbb{T}

$$\{x_{\alpha}, x_{\beta}\} = \langle \alpha, \beta \rangle \cdot x_{\alpha+\beta}.$$

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$$\text{DT}_\sigma(\gamma) = "e(\mathcal{M}^{\sigma-ss}(\gamma))" \in \mathbb{Q}, \quad \gamma \in \Gamma,$$

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There are equivalent invariants $\Omega_\sigma(\gamma)$ defined by

$$\text{DT}_\sigma(\alpha) = \sum_{\alpha=k\cdot\beta} \frac{1}{k^2} \cdot \Omega_\sigma(\beta).$$

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$$\mathbb{S}(\ell)^*(x_\beta) = \exp \left\{ \sum_{Z(\gamma) \in \ell} \text{DT}(\gamma) \cdot x_\gamma, - \right\} (x_\beta) = x_\beta \prod_{Z(\gamma) \in \ell} (1 - x_\gamma)^{\Omega(\gamma) \cdot \langle \beta, \gamma \rangle}$$

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- use automorphisms defined on analytic open subsets of \mathbb{T} .

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Controls the variation of the DT invariants as $\sigma \in \text{Stab}(\mathcal{D})$ varies.

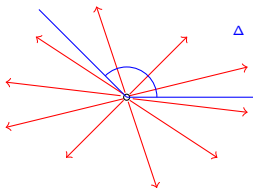
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For any convex sector $\Delta \subset \mathbb{C}^*$, the clockwise ordered product

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is constant, providing no Stokes ray crosses $\partial\Delta$.



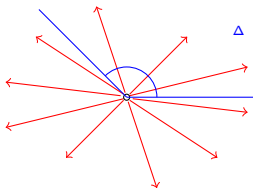
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The Stokes rays are spanned by the points $Z(\gamma)$ with $\text{DT}(\gamma) \neq 0$; the formula makes sense in the completed torus algebra.

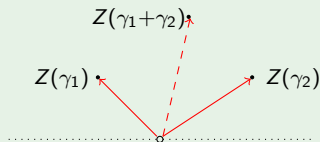
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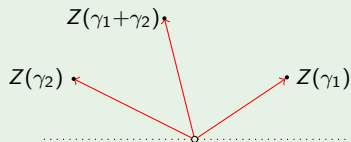
Here $\Gamma = \mathbb{Z}\gamma_1 \oplus \mathbb{Z}\gamma_2$ with $\langle \gamma_1, \gamma_2 \rangle = 1$, and $\mathbb{C}[\mathbb{T}] = \mathbb{C}[x_1^{\pm 1}, x_2^{\pm 1}]$.

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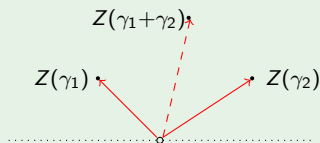
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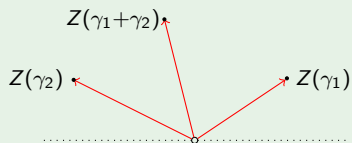
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The wall-crossing formula is the cluster pentagon identity

$$C_{\gamma_1} \circ C_{\gamma_2} = C_{\gamma_2} \circ C_{\gamma_1 + \gamma_2} \circ C_{\gamma_1},$$

$$C_{\alpha} : x_{\beta} \mapsto x_{\beta} \cdot (1 + x_{\alpha})^{\langle \alpha, \beta \rangle}.$$

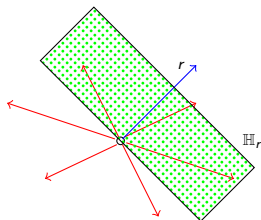
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Our stability condition gives rise to a collection of rays in \mathbb{C}^* labelled by elements $\mathbb{S}(\ell)$ of the (slightly mythical) group $G = \text{Aut}(\mathbb{T})$.

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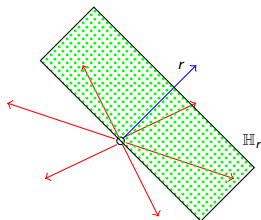
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This is motivated by the wall-crossing formula, and an analogy with Stokes data of differential equations, as appearing in Dubrovin's work on semisimple Frobenius manifolds.

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and if $\Delta \subset \mathbb{C}^*$ is a convex sector with $\partial\Delta = \{r_+\} \cup \{r_-\}$ then

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We are taking $G = \text{Aut}_{\{-, -\}}(\mathbb{T})$, and composing $Y: \mathbb{C}^* \rightarrow G$ with the evaluation map $\text{ev}_\xi: G \rightarrow \mathbb{T}$.

4. Example: the A_1 case.

See: "Riemann-Hilbert problems from Donaldson-Thomas theory", arxiv.

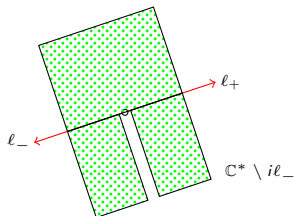
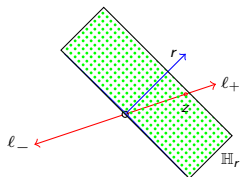
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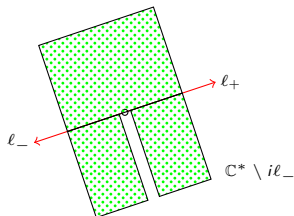
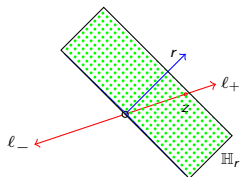
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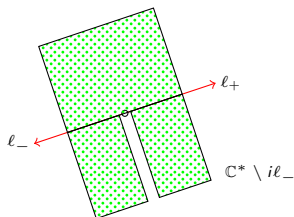
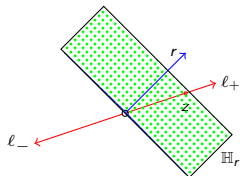
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$$X_-(t) = X_+(t) \quad \hat{X}_-(t) = \begin{cases} \hat{X}_-(t) \cdot (1 - X(t)^{+1}) & \operatorname{Re}(t/z) > 0 \\ \hat{X}_-(t) \cdot (1 - X(t)^{-1}) & \operatorname{Re}(t/z) < 0 \end{cases}$$

As $t \rightarrow 0$:

$$X_{\pm}(t) \cdot e^{z/t} \rightarrow 1 \quad \text{and} \quad \hat{X}_{\pm}(t) \cdot e^{\hat{z}/t} \rightarrow 1,$$

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The unique solution to the problem is

$$X_{\pm}(t) = e^{-z/t}, \quad \hat{X}_{\pm}(t) = e^{-\hat{z}/t} \cdot \Lambda\left(\frac{\pm z}{2\pi i t}\right)^{\pm 1},$$

where $\Lambda(s)$ is the modified gamma function

$$\Lambda(s) = \frac{e^s \cdot \Gamma(s)}{\sqrt{2\pi} \cdot s^{s-\frac{1}{2}}} \sim \exp\left(\sum_{g=1}^{\infty} \frac{B_{2g} \cdot s^{1-2g}}{2g(2g-1)}\right).$$

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Then we can write $\tau(z, t) = \Upsilon(z/2\pi it)$ with

$$\Upsilon(w) = \frac{e^{-\zeta'(-1)} \cdot e^{\frac{3}{4}w^2} \cdot G(w+1)}{(2\pi)^{\frac{w}{2}} \cdot w^{\frac{w^2}{2}}} \sim \sum_{g \geq 2} \frac{B_{2g} \cdot w^{2-2g}}{2g \cdot (2g-2)}$$

a modified Barnes G -function.

5. Geometric case: coherent sheaves on a CY_3 .

See: "Riemann-Hilbert problems and the resolved conifold", arxiv.

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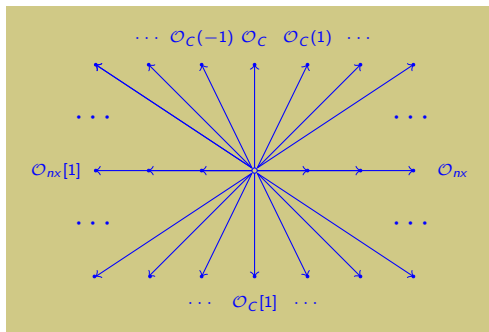
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The space of stability conditions is a cover of

$$M = \{(v, w) \in \mathbb{C}^2 : w \neq 0 \text{ and } v + dw \neq 0 \text{ for all } d \in \mathbb{Z}\} \subset \mathbb{C}^2$$

with central charge $Z(r, d) = rv + dw$.

DT INVARIANTS AND RAY DIAGRAM



$$\Omega(\gamma) = \begin{cases} 1 & \text{if } \gamma = \pm(1, d) \text{ for some } d \in \mathbb{Z}, \\ -2 & \text{if } \gamma = (0, d) \text{ for some } 0 \neq d \in \mathbb{Z}, \\ 0 & \text{otherwise.} \end{cases}$$

NON-PERTURBATIVE PARTITION FUNCTION

There is a unique solution to the RH problem which can be written explicitly using Barnes double and triple gamma functions.

$$H(v, w, t) = \int_{\mathbb{R}+i\epsilon} \frac{e^{vs} - 1}{e^{ws} - 1} \cdot \frac{-e^{ts}}{(e^{ts} - 1)^2} \cdot \frac{ds}{s}$$

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$$\begin{aligned} \tau(v, w, t) &= \exp(H + R) \sim \text{sporadic } g = 0, 1 \text{ terms} + \\ &+ \sum_{g \geq 2} \frac{B_{2g}}{2g(2g-2)!} \left(\text{Li}_{3-2g}(e^{2\pi iv/w}) + \frac{B_{2g-2}}{2g-2} \right) \left(\frac{2\pi it}{w}\right)^{2g-2} \end{aligned}$$

5. Further directions.

COMPLEX HYPERKÄHLER STRUCTURE

Some joint work with Ian Strachan.

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Can we construct the twistor space using just the DT invariants?
Then the DT RH problem is just constructing the twistor lines?

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Consider $\mathcal{D} \subset \mathcal{D}^b \text{Lag}(X, \omega)$ with X locally of the form

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There is a double cover $\hat{S} \rightarrow S$ defined by $y^2 = q(x)$ and

$$\mathbb{T} = H^1(\hat{S}, \mathbb{C}^*)^-$$

COMPLEXIFIED HITCHIN SYSTEM

Solve the RH problems geometrically by studying a moduli space

$$\mathcal{M} = \{(S, E, \nabla, \Phi)\},$$

of a Riemann surface S , a rank 2 vector bundle E on S , equipped with both a holomorphic connection ∇ and a Higgs field Φ .

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The spectral curve construction shows that the map

$$\pi: \mathcal{M} \rightarrow \mathcal{Q}, \quad (S, E, \nabla, \Phi) \mapsto (S, \det(\Phi)),$$

has fibres which are copies of \mathbb{T} .